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# High-intensity spatial-mode steerable frequency up-converter toward on-chip integration

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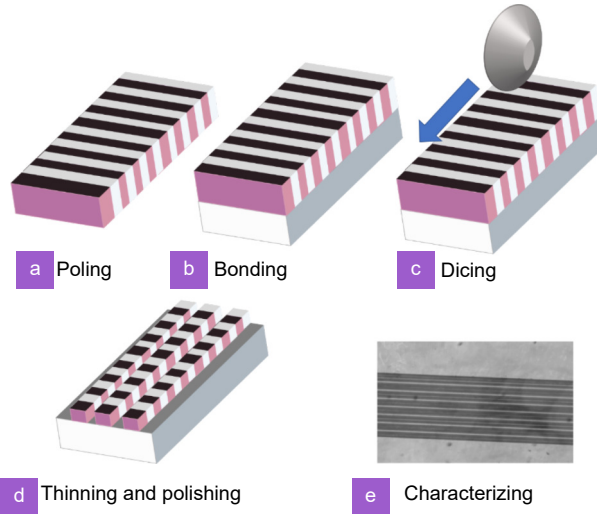
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**Section 1: Spatial modes of the ridge waveguide**

For the Z-cut PPMgLN waveguide in Fig. 1(a), spatial modes inside the waveguide were calculated using COMSOL Multiphysics Software using the frequency domain module for electromagnetic waves. The effective indices of the PPLN waveguide at a broad temperature range can be calculated synchronously for the pump, signal, and SFG lights. Here, temperature-dependent Sellmeier formula of the LN crystal was used <sup>S1</sup>:



**Fig. S1 | Flowchart for preparing the spatial-mode steerable SFG waveguide.**

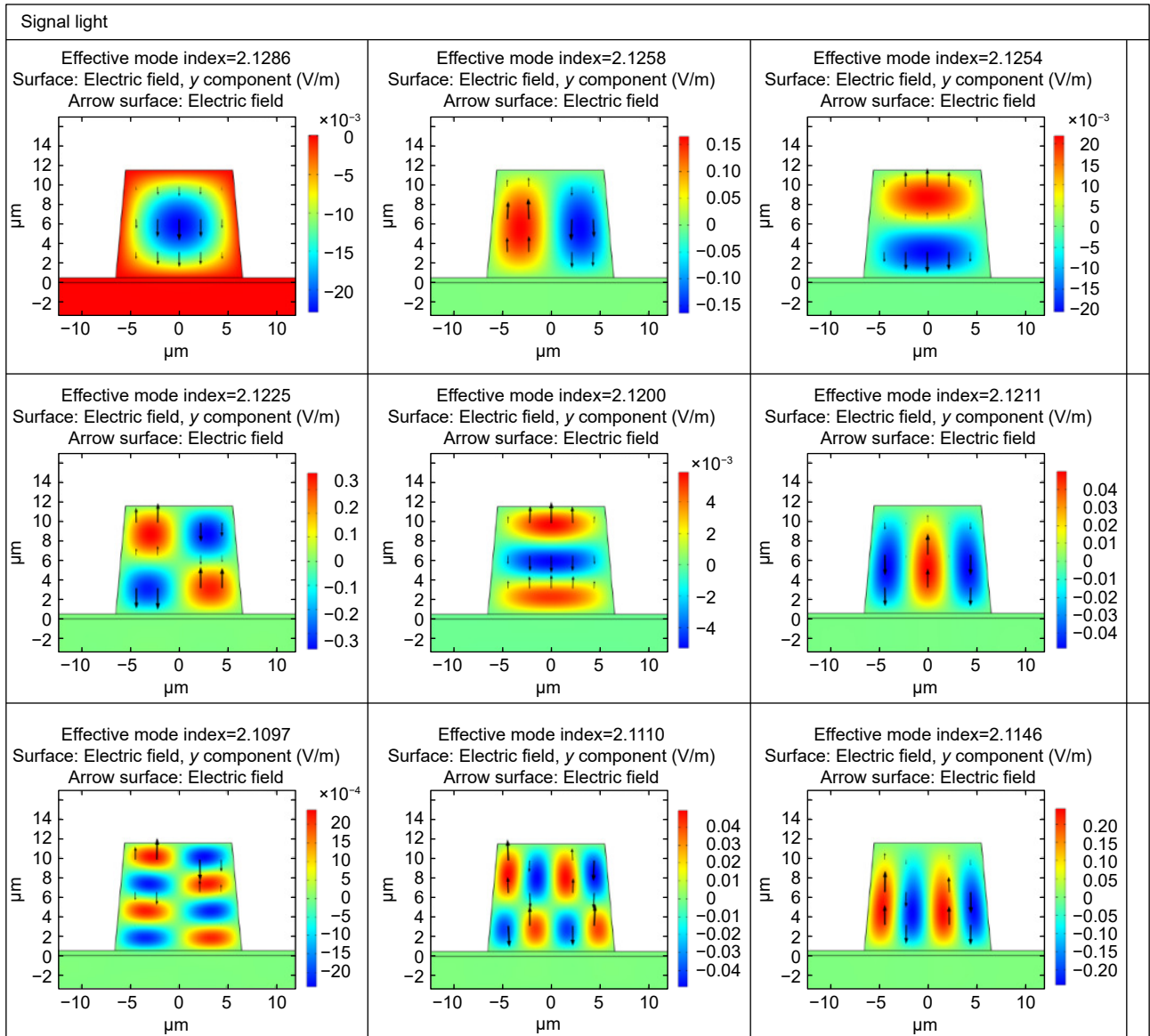
$$n_e(T) = \sqrt{A_{e1} + fB_{e1} + \frac{A_{e2} + fB_{e2}}{\lambda^2 - (A_{e3} + fB_{e3})^2} + \frac{A_{e4} + fB_{e4}}{\lambda^2 - A_{e5}^2} - A_{e6}\lambda^2}, \tag{S1}$$

where  $f(T) = (T-24.5)(T+570.82)$ , and parameters from  $A_{e1}$  to  $A_{e6}$  and  $B_{e1}$  to  $B_{e2}$  are listed in Table S1.

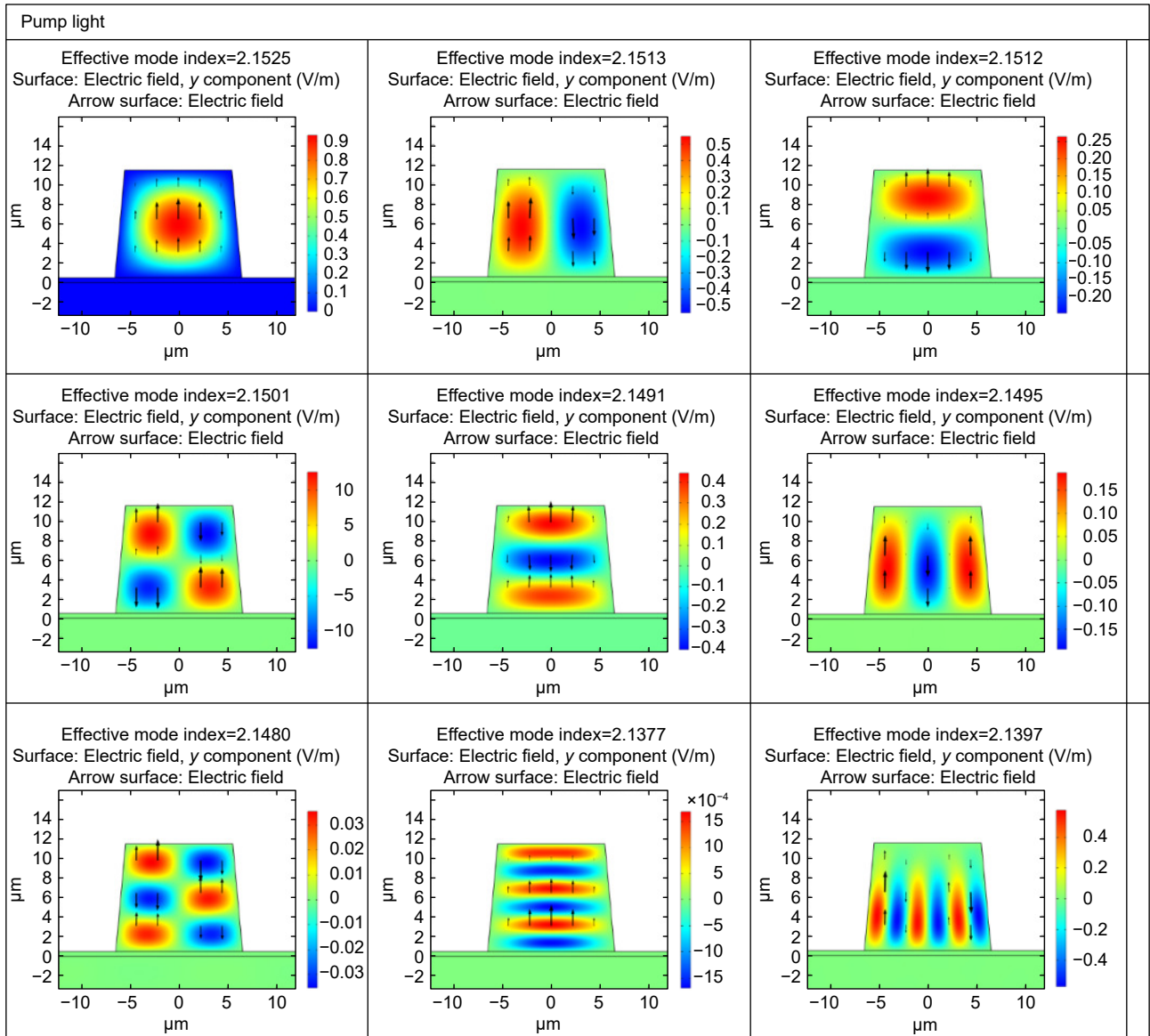
**Table S1 | Sellmeier parameters for the LN crystal**

$A_{e1}$	$A_{e2}$	$A_{e3}$	$A_{e4}$	$A_{e5}$	$A_{e6}$
5.756	0.0983	0.2020	189.32	12.52	$1.32 \times 10^{-2}$
$B_{e1}$	$B_{e2}$	$B_{e3}$	$B_{e4}$		
$2.86 \times 10^{-6}$	$4.7 \times 10^{-8}$	$6.113 \times 10^{-8}$	$1.516 \times 10^{-4}$		

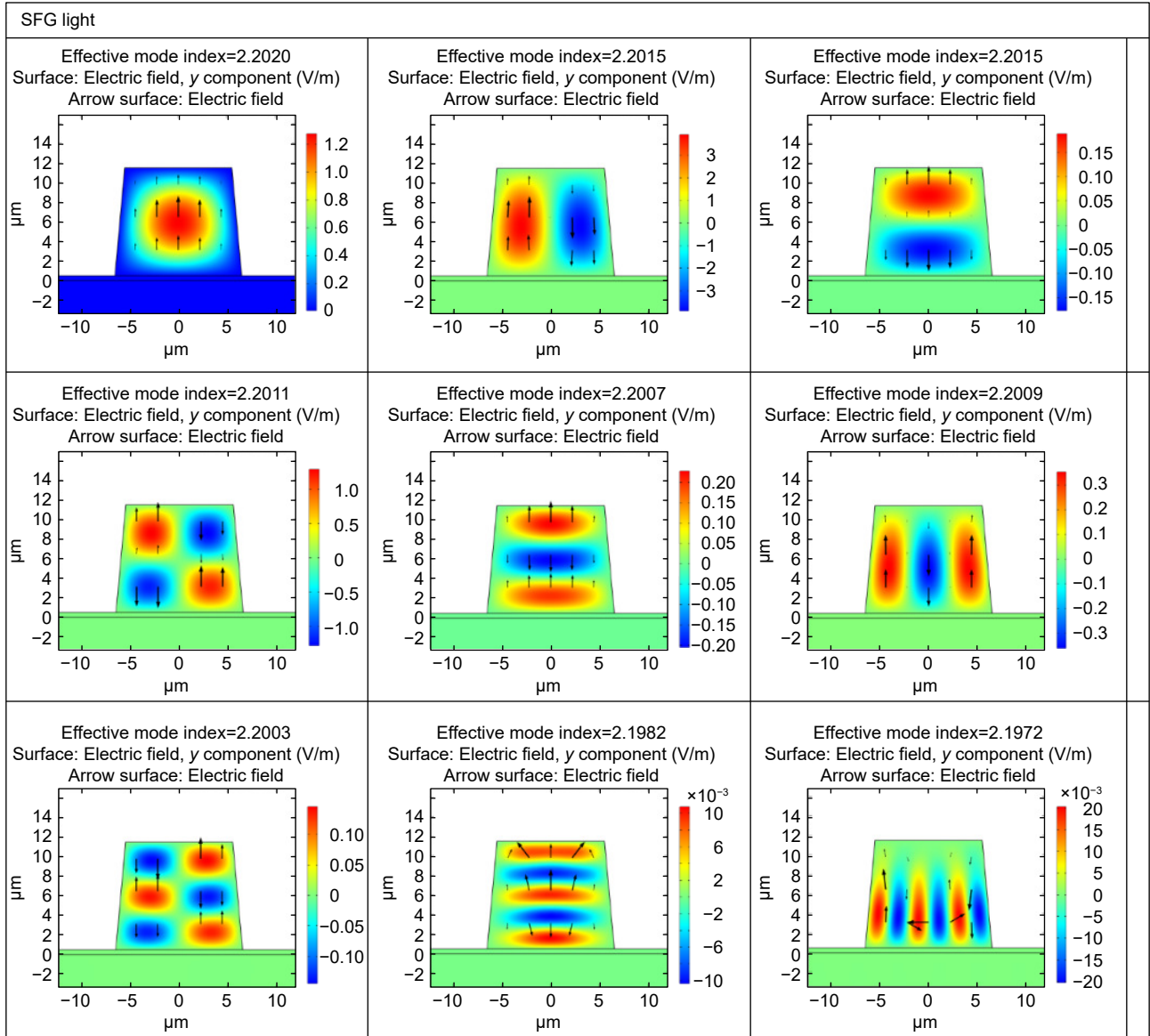
**Table S2 | Typical spatial modes among the 11 TM modes of the 1552.6-nm signal light at 25°C**



**Table S3 | Typical spatial modes among the calculated 47 TM modes of the 973.85-nm pump light at 25°C**



**Table S4 | Typical spatial modes among the calculated 50 TM modes of the 598.47-nm SFG light at 25°C**



## Section 2: Calculating the inter-mode quasi-phase matching efficiency

Combination of the guided signal, pump, and SFG lights can be defined as follows:

$$A_i(x, y, z) = \sum_{\mu\nu} A_i^{\mu\nu}(z) \varphi_i^{\mu\nu}(x, y) \exp(-i\beta_i^{\mu\nu} z), \quad (S2)$$

where  $i = P, S,$  and  $F$  denote the pump, signal, and SFG lights, respectively, and  $\varphi^{i\mu\nu}$  defines the spatial profile of mode  $i\mu\nu$ . Since the signal and pump lights are with high incident powers which are approximately un-depleted during SFG<sup>S2</sup>, the amplitude of the SFG light at existing plane of the waveguide is expressed as follows:

$$A_F(x, y, L) = \frac{4\pi L i d_{\text{eff}} A_P A_S}{n_F \lambda_F} \left( \frac{\exp(i\Delta k L) - 1}{i\Delta k L} \right), \quad (S3)$$

For  $I_i = 2n_i \epsilon_0 c |A_i|^2$ , the SFG intensity is as follows:

$$I_F(x, y, L) = \frac{8\pi^2 L^2 d_{\text{eff}}^2 I_P I_S}{n_F n_S n_P \epsilon_0 c \lambda_F^2} \text{sinc}^2 \left( \frac{\Delta k L}{2} \right), \quad (S4)$$

The power of the SFG light could be calculated by integrating Eq. S4 over the waveguide cross-section. Considering

the inter-mode conversation, overlap integral  $\Phi$  among modes  $\lambda_S^{ml}$ ,  $\lambda_P^{uv}$ , and  $\lambda_F^{jk}$  at temperature  $T$  was introduced:

$$\Phi(x, y, T) = \frac{\left| \iint E_S^{ml}(x, y, T) E_P^{uv}(x, y, T) E_F^{jk*}(x, y, T) dx dy \right|^2}{\iint |E_S^{ml}(x, y, T)|^2 dx dy \iint |E_P^{uv}(x, y, T)|^2 dx dy \iint |E_F^{jk*}(x, y, T)|^2 dx dy} \quad (S5)$$

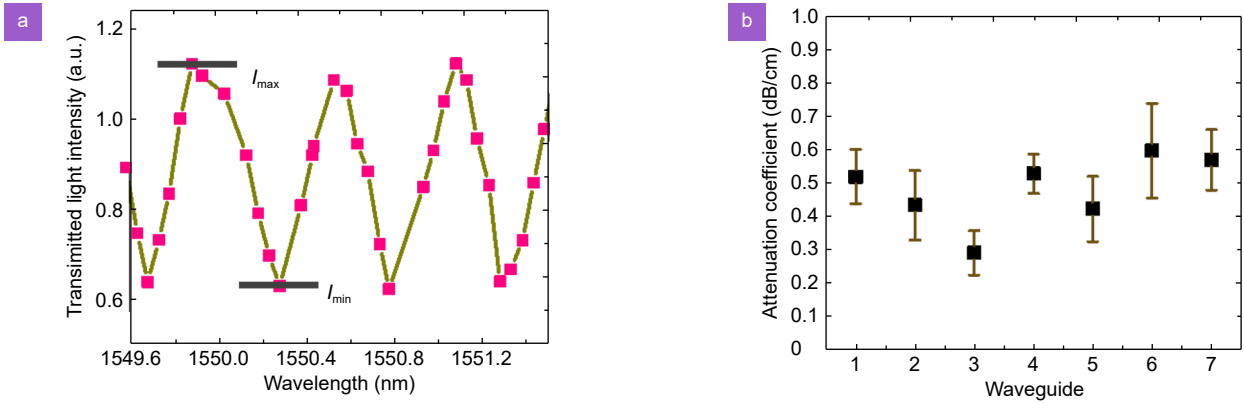
Here,  $E_S^{ml}$ ,  $E_P^{uv}$ , and  $E_F^{jk}$  denote the amplitudes of the electric field y component (TM component) for modes  $\lambda_S^{ml}$ ,  $\lambda_P^{uv}$ , and  $\lambda_F^{jk}$ , respectively, at temperature  $T$ , and the SFG power for  $\lambda_F^{jk}$  can be revised as follows:

$$P_F^{jk}(x, y, T) = \frac{8\pi^2 L^2 d_{\text{eff}}^2 P_P^{uv} P_S^{ml}}{n_S^{ml} n_P^{uv} n_F^{jk} \epsilon_0 c \lambda_F^2} \Phi(x, y, T) \text{sinc}^2\left(\frac{\Delta k L}{2}\right) \quad (S6)$$

### Section 3: Device fabrication and characterization

First, the poling period was fabricated in a 500- $\mu\text{m}$ -thick Z-cut MgO: LN wafer using the standard applied pulsed voltage polarization technique (Fig. S1(a)). Next, the PPLN crystal was bonded to a 500- $\mu\text{m}$ -thick LT substrate, with the assistance of a 0.5- $\mu\text{m}$ -thick SiO<sub>2</sub> layer (Fig. S1(b)). A diamond saw was used to dice and remove the additional material (Fig. S1(c)) to form a 30-mm long PPLN waveguide array consisting of seven waveguides (Fig. S1(f)), after thinning and fine polishing (Fig. S1(d)).

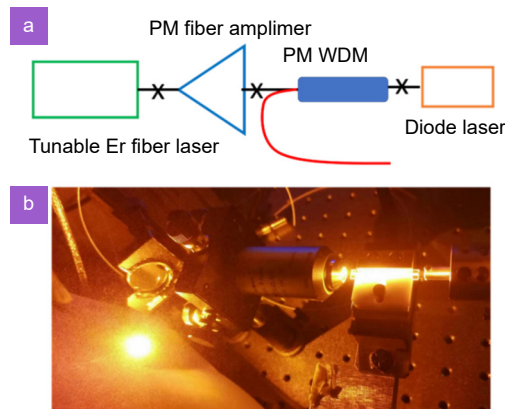
A 3D profiler (S neox, Sensofar Inc., Spain) was used to characterize the detailed structure of the waveguide array and the poling period. The attenuation coefficient of each waveguide was detected by measuring the contrast of Fabry–Perot resonances with a tunable lasing system at the C-band (Fig. S2), where the third waveguide with the lowest loss of 0.29 dB/cm was used for demonstrating spatial mode steering experiments.



**Fig. S2 | Attenuation loss of the prepared waveguide array.** (a) Measured transmitted intensity by varying incident wavelength of the C-band laser to the third waveguide. (b) Calculated attenuation coefficient for each waveguide according to the method in S3, S4. The following relationships are used: the attenuation coefficient  $\alpha = 4.34(\ln R - \ln R')/L$ , the end-face reflectivity  $R = (n_{\text{eff}} - 1)^2 / (n_{\text{eff}} + 1)^2$ , the combined loss-reflection factor  $R' = [1 - (1 - K^2)^{1/2}] / K$ , and the contrast of the Fabry–Perot resonances  $K = (I_{\max} - I_{\min}) / (I_{\max} + I_{\min})$ . Here,  $n_{\text{eff}} = 2.212$  is the effective index at 1550 nm,  $L = 30$  mm is the waveguide length, and  $I_{\max}$  and  $I_{\min}$  denote the peak and valley of the transmitted light intensity in a, respectively.

## Section 4: Experimental details

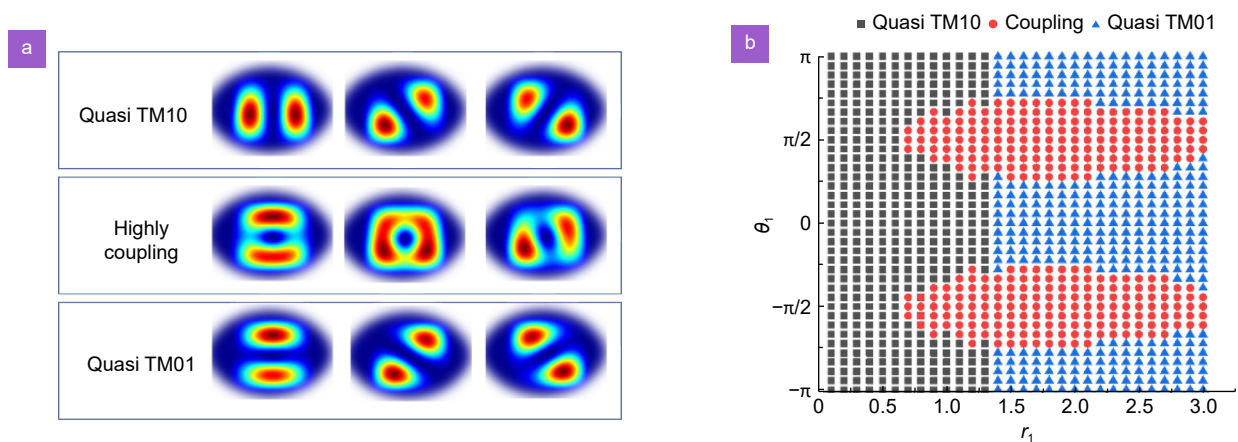
A tunable linearly polarized Er-doped fiber laser at C-band (YO13040139, Yenista Inc., French), with a linewidth of 17 pm and a tuning range from 1546 to 1556 nm, was used to provide the signal wavelength, which was amplified by an all-fiber amplification stage to increase the signal power from 10 to 27.79 dBm (Fig. S3(a)). A fiber-coupled polarization-maintaining single-mode LD, with a maximum output power of 26.02 dBm at 973.85 nm was used as the pump source. The signal and pump lights were coupled with a wavelength division multiplexer before entering the fabricated few-mode PPLN waveguide. The layout of the light sources for SFG is displayed in Fig. 2(a). A bandpass filter was placed after the waveguide for filtering SFG light at 598 nm. The SFG power was measured with a power meter (3 A, Ophir Optonics Solutions, Ltd., Israel); the pump and SFG spectra by a spectral analyzer, with a spectral resolution of 0.02 nm (AQ6370D, Yokogawa Electric Corporation, Japan). To catch the spatial mode profile, a 20X object lens (Sigma Inc., Japan) was placed in front of the waveguide output surface (Fig. S3(b)). The spatial mode profile was measured using a CCD camera (COMS-1202, Cinogy Inc., Germany).



**Fig. S3** | (a) Layout of the light source consisting of pump and signal lights; (b) image of the spatial mode steerable SFG experiment.

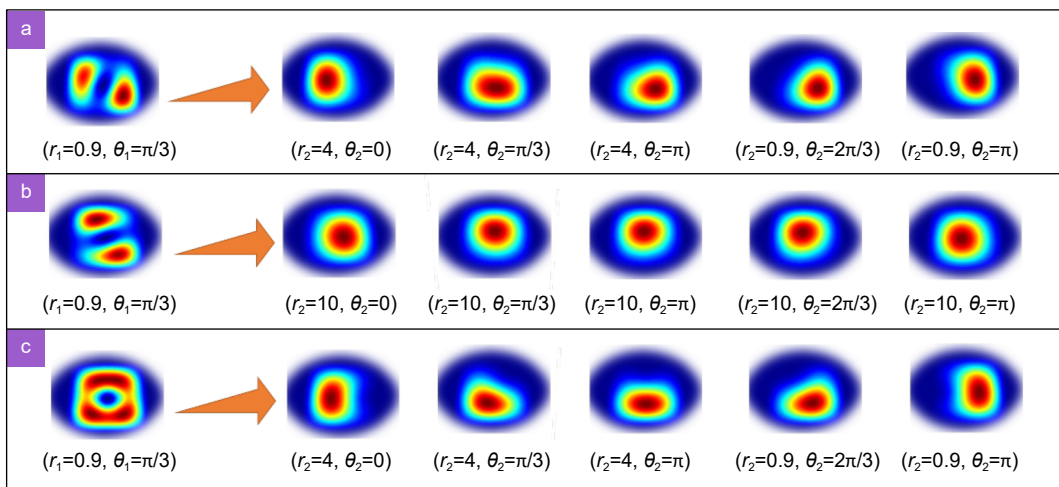
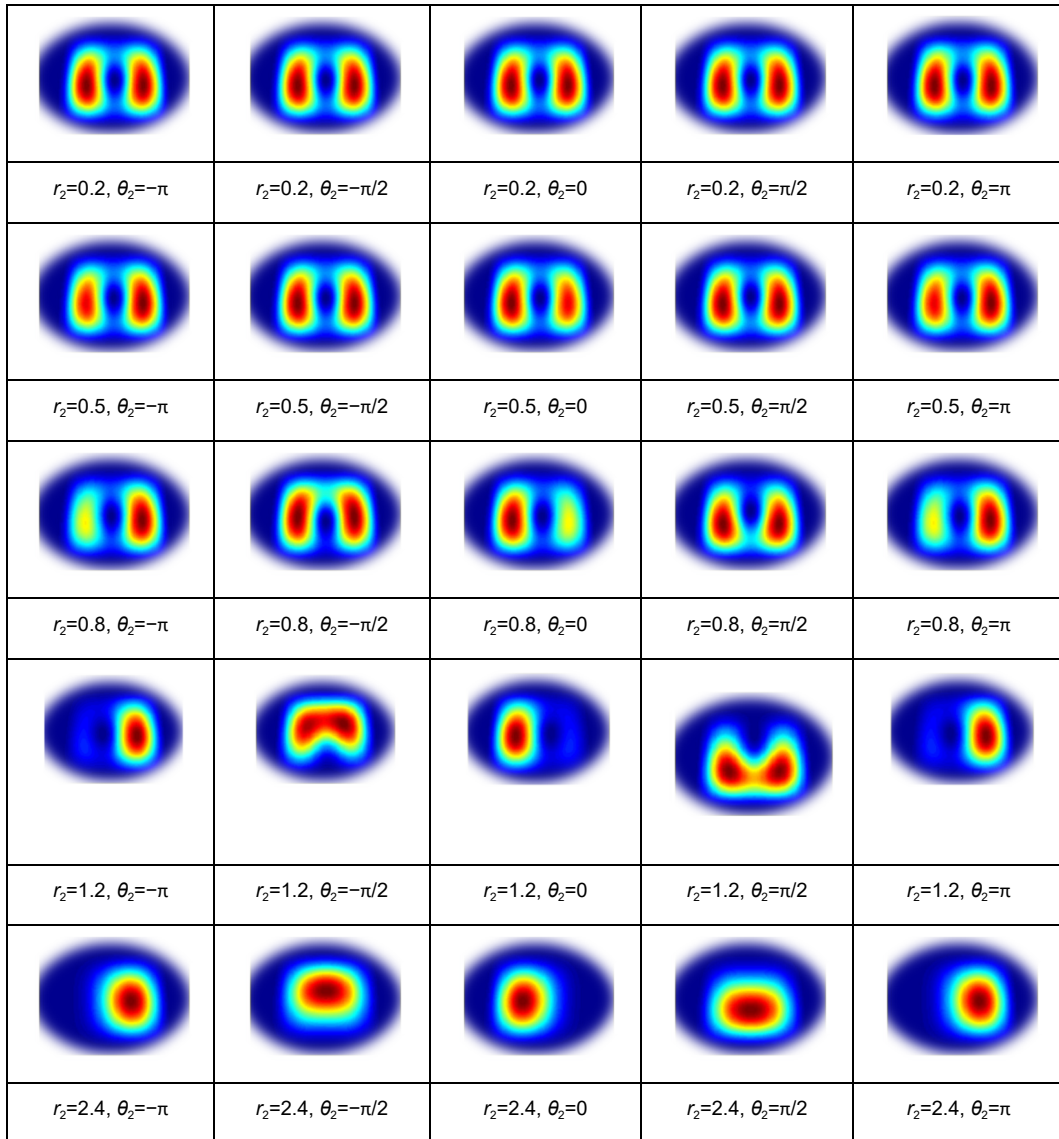
## Section 5: Evaluating the coupling modes

In the two-mode coupling state between TM01 and TM10, we obtained the coupling mode matrix by varying  $r_1$  from 0.1 to 3 with a step of 0.1 and varying  $\theta_1$  from  $-\pi$  to  $\pi$  with a step of  $\pi/180$ . By matching the mode pattern in the matrix with a given coupling mode, that is, the patterns in Fig. 4(d), the relative amplitude and phase angle between TM01 and TM10 can be estimated. In the three-mode coupling state among TM01, TM10, and TM00, a primary pattern in the two-mode coupling matrix was selected to match the three-mode coupling pattern before sweeping the parameters between  $r_2$  and  $\theta_2$  to simulate the three-mode coupling pattern (Table S5).



**Fig. S4** | (a) Mode patterns classified as the quasi TM10, quasi TM01, and the high coupling modes; and (b) parameter range for the high coupling mode between TM10 and TM01.

**Table S5 | Simulating the three-mode coupling process**

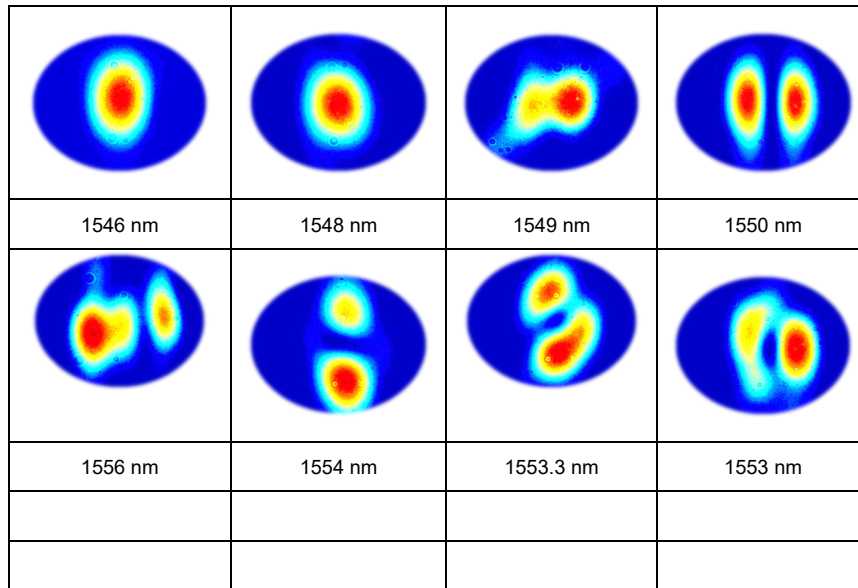


**Fig. S5 |** Three-mode coupling pattern under sufficiently high TM00 mode intensity for the initial (a) quasi TM01, (b) quasi TM10, and (c) highly coupling modes defined in Fig. S4.

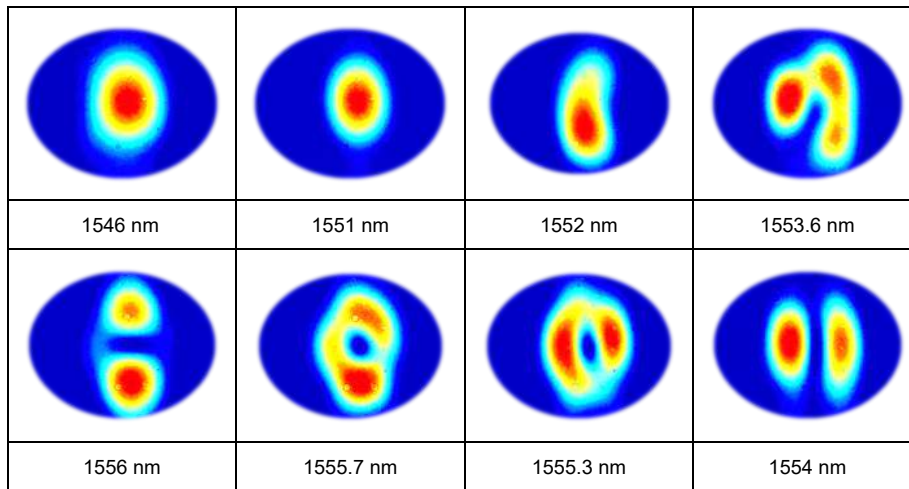


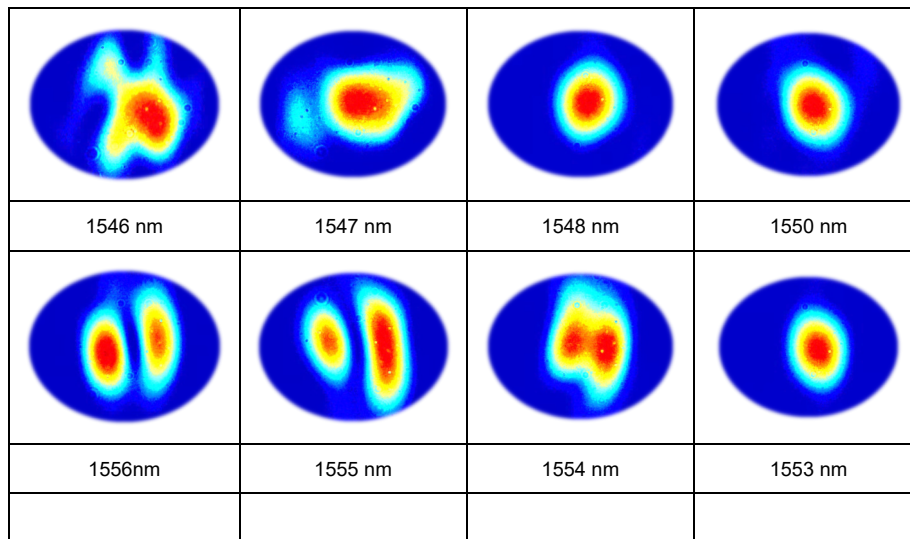
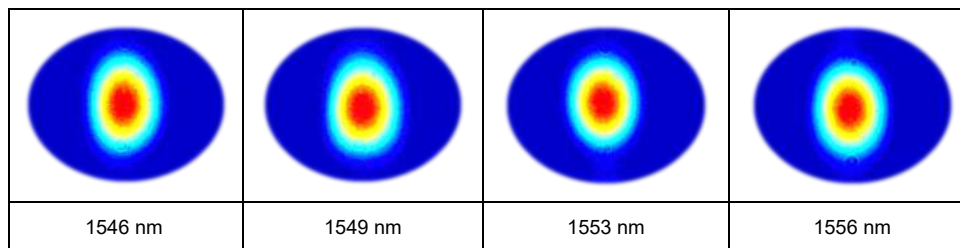
Section 6: Wavelength-dependent spatial mode steering scheme

**Table S6 | Evolution of the spatial mode in the wavelength-dependent spatial mode steering SFG scheme at 35°C**



**Table S7 | Evolution of the spatial mode in the wavelength-dependent spatial mode steering SFG scheme at 45°C**



**Table S8 | Evolution of the spatial mode in the wavelength-dependent spatial mode steering SFG scheme at 55°C****Table S9 | Evolution of the spatial mode in the wavelength-dependent spatial-mode steering SFG scheme at 65°C**

## References

- S1. Gayer O, Sacks Z, Galun E, Arie A. Temperature and wavelength dependent refractive index equations for MgO-doped congruent and stoichiometric LiNbO<sub>3</sub>. *Appl Phys B* **91**, 343–348 (2008).
- S2. Boyd R. *Nonlinear Optics* (Academic Press, 2008).
- S3. Regener R, Sohler W. Loss in low-finesse Ti: LiNbO<sub>3</sub> optical waveguide resonators. *Appl Phys B* **36**, 143–147 (1985).
- S4. Dong P, Qian W, Liao SR, Liang H, Kung CC et al. Low loss shallow-ridge silicon waveguides. *Opt Express* **18**, 14474–14479 (2010).